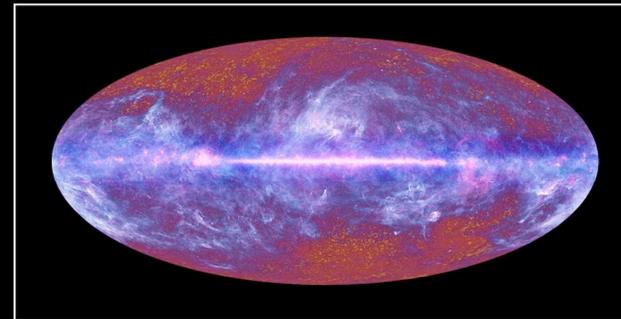
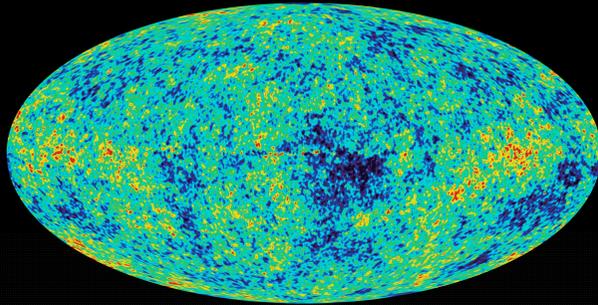


**The MOND phenomenology and its
covariant theories:
successes and failures**

B. Famaey

The Λ CDM model of the Universe

(i) Cosmic Microwave Background (WMAP, Planck soon)



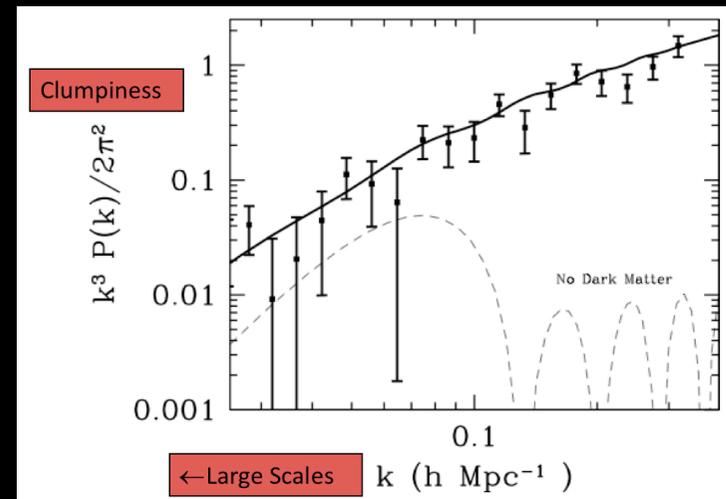
The Planck one-year all-sky survey



(c) ESA, HFI and LFI consortia, July 2010

(ii) Matter power spectrum (SDSS)

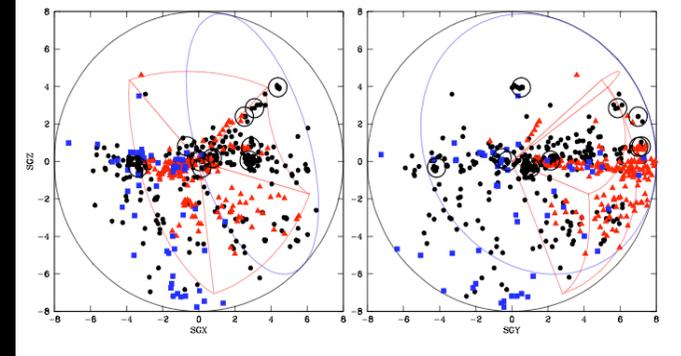
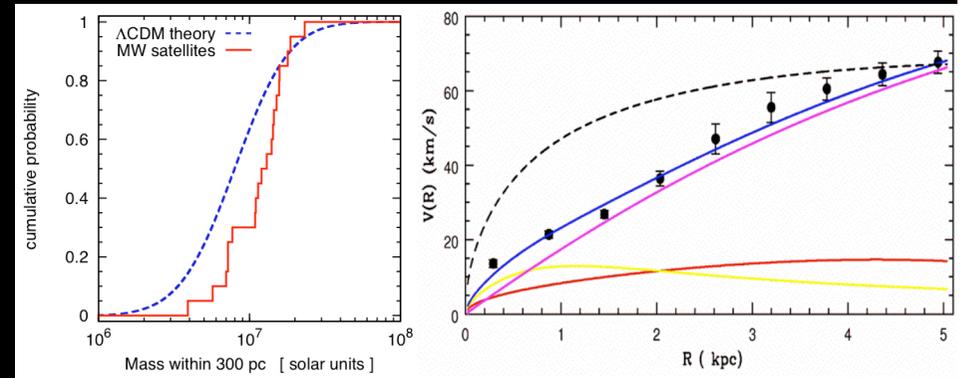
(iii) Type Ia supernovae



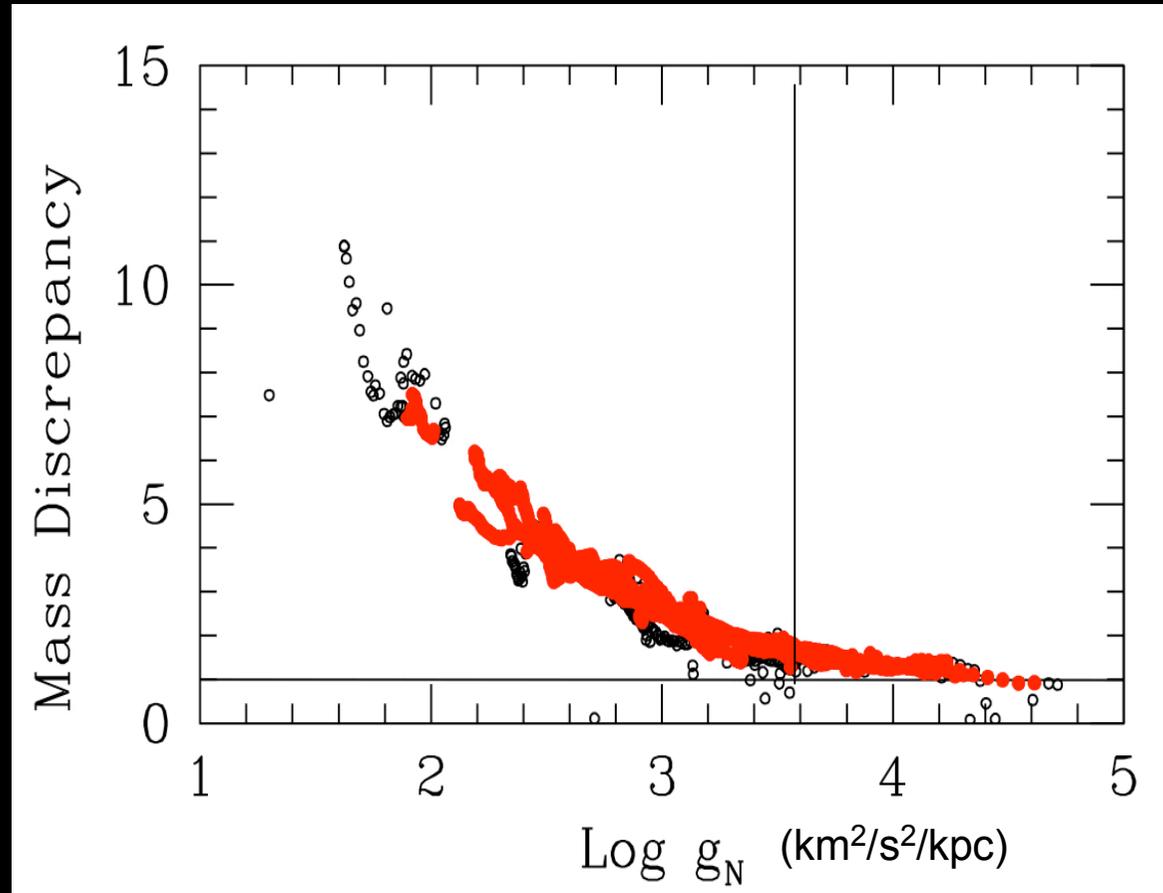
(iv) Weak lensing; (v) strong lensing; (vi) temperature profiles of intracluster gas; (vii) velocity dispersion profiles of pressure-supported galaxies; (viii) rotation curves of spiral galaxies (?)

Challenges

- (i) Fine-tuned value of Λ
- (ii) Missing satellite problem
- (iii) Cusp problem
- (iv) Angular momentum problem
- (v) Phase-space correlation problem
- (vi) Local Void problem
- (vii) Scaling relations :
 - (i) Tully-Fisher $V_{\infty}^4 \propto M_{\text{bar}}$: slope, zero point and small scatter;
 - (ii) Mean surface density within r_0 ;
 - (iii) one-to-one correlation between total and baryonic gravitational fields (MDA or the « conspiracy » problem)



Mass Discrepancy vs Acceleration



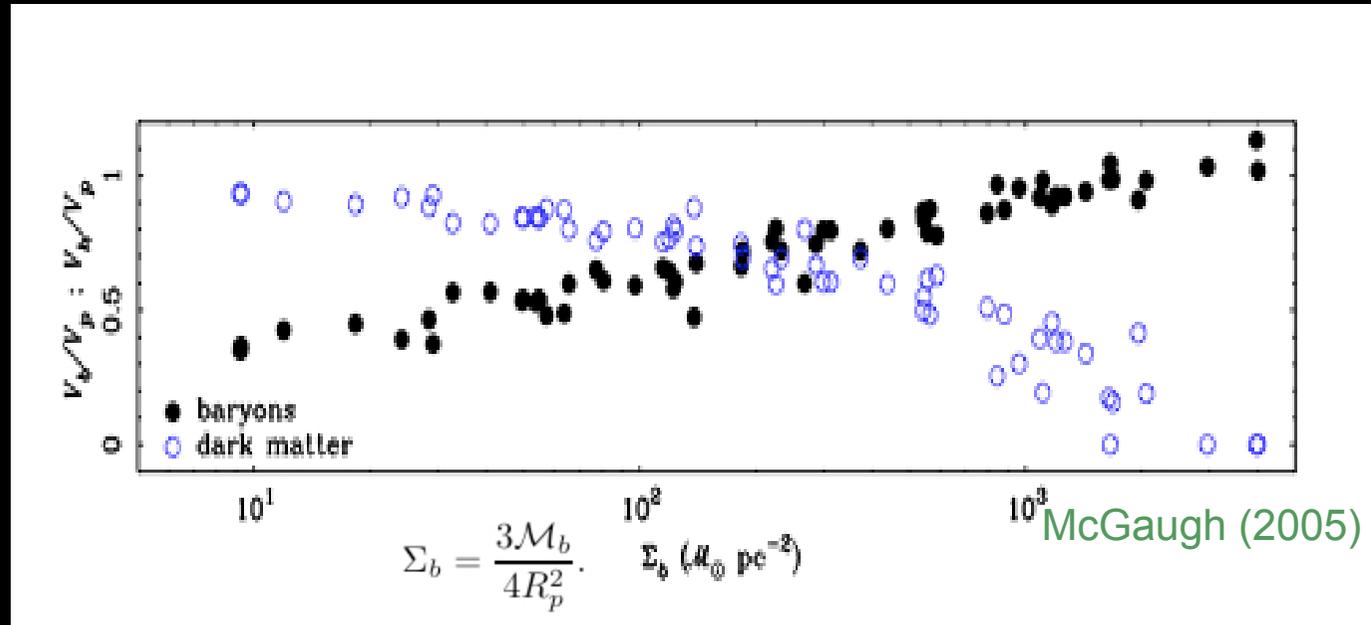
At small a :

$$M_{\text{tot}}/M_{\text{bar}} = a^{-1}$$

in units of the
transition
acceleration a_0
such that
 $\Lambda = (a_0/c)^2$
in s^{-2}

McGaugh 2004; Gentile, Famaey & de Blok 2010

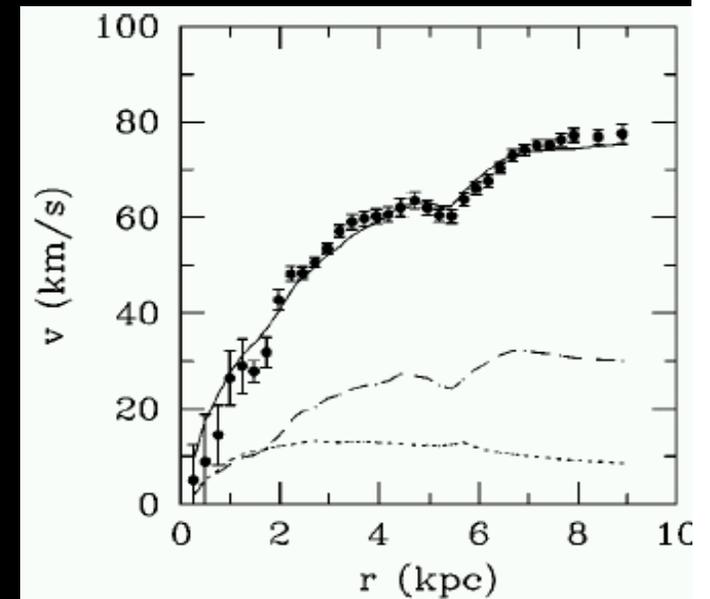
A fine balance of DM and baryons



- R_p = radius of max contribution of both gas and stars to the RC
- Comparing the contribution of baryons to the RC as a function of surface density (proxy for characteristic acceleration)
- This could point at some **repulsion** between surface densities of baryons and DM
- MDA (and this MDsurfdn) is *history-independent* !

Sancisi's rule: locally no repulsion

Each time one sees a feature in the light, there is a feature in the rotation curve (Sancisi's rule)



⇒ Dark baryons as dark matter?

(Pfenniger & Combes 1994)

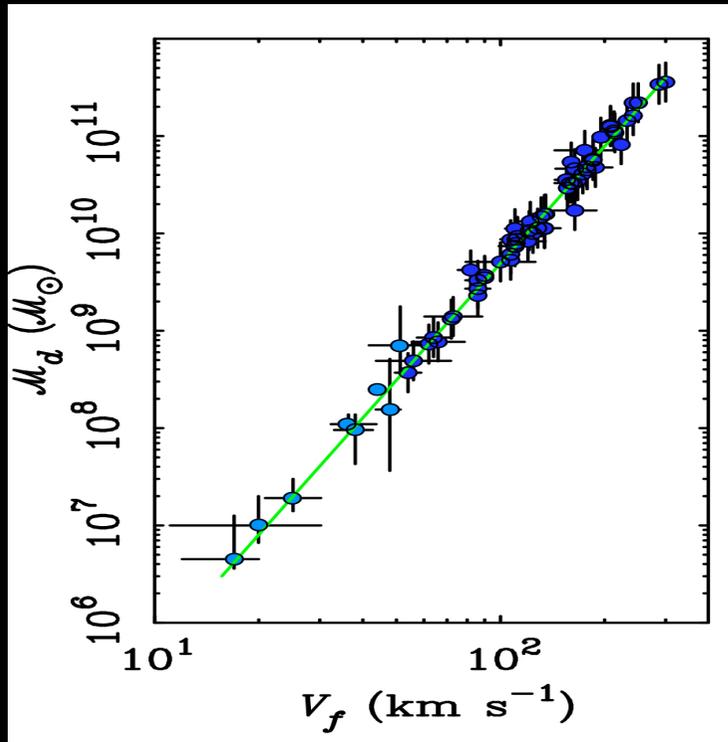
MDA asymptotes to BTF

At small accelerations, the mass discrepancy

$$\text{is } \boxed{M_{\text{tot}}/M_{\text{bar}} = a_0/a}$$

$$\Rightarrow V^2 = GM_{\text{tot}}/r = GM_{\text{bar}}a_0/(ar) = GM_{\text{bar}}a_0/V^2$$

$$\Rightarrow V^4 \propto M_{\text{bar}}$$



Evolution of the BTF with z ?

If the BTF evolves or not, and if there is something fundamental about the value of a_0 , it would allow to know whether the fundamental relation is :

$$a_0/c = H_0$$

or

$$(a_0/c)^2 = \Lambda$$

or

a coincidence

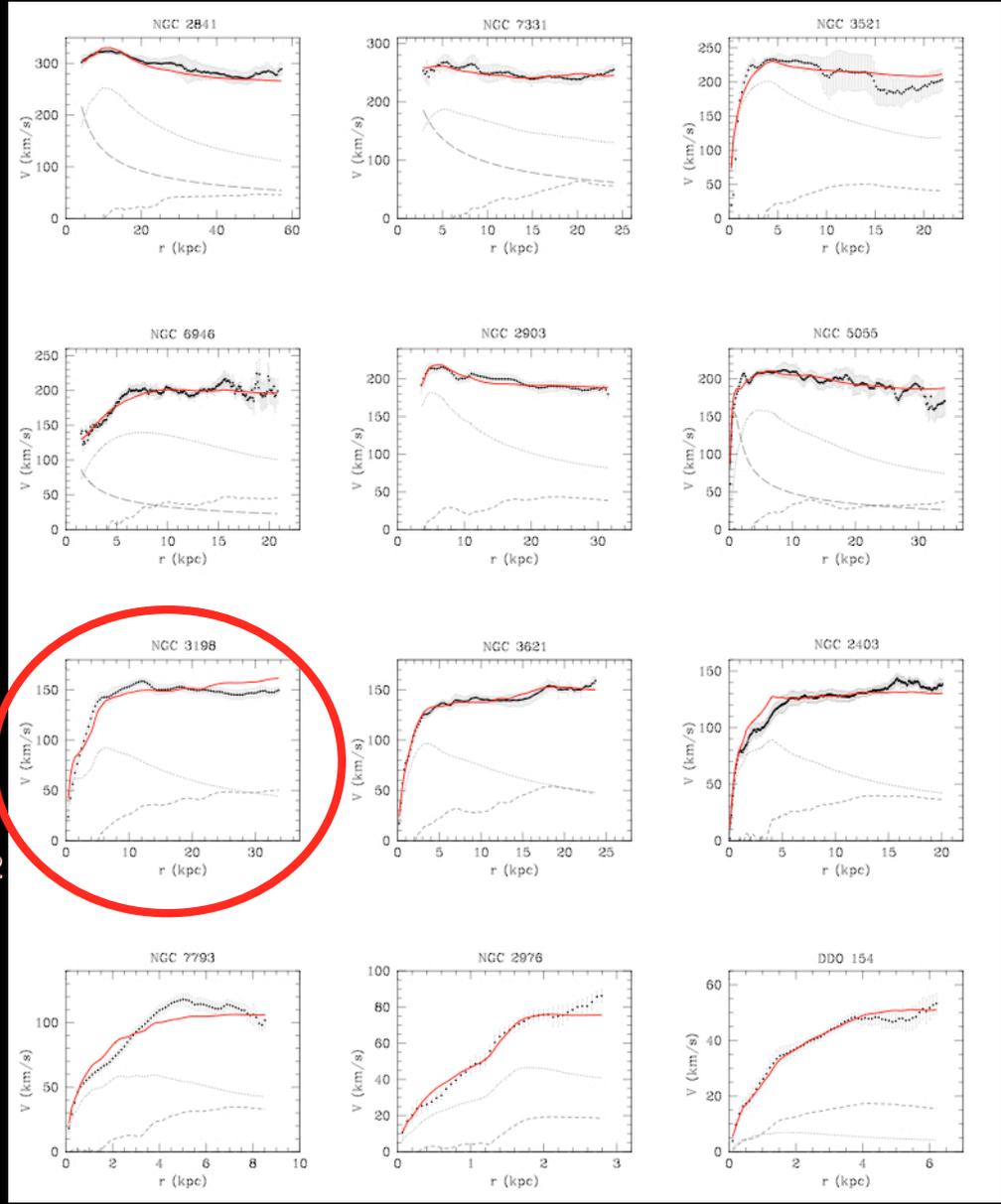
The MDA can be summarized by Milgrom's formula

Correlation summarized by this formula in galaxies
(Milgrom 1983):

$$\begin{aligned} \mu(g/a_0) g &= g_{\text{N bar}} & \text{or} & & \nu(g_{\text{N bar}}/a_0) g_{\text{N bar}} &= g \\ \text{with } \mu(x) &= x & \text{or } \nu(x) &= x^{-1/2} & \text{for } x \ll 1 & \text{(deep-MOND)} \\ \mu(x) &= \nu(x) = 1 & & & \text{for } x \gg 1 & \text{(Newtonian)} \end{aligned}$$

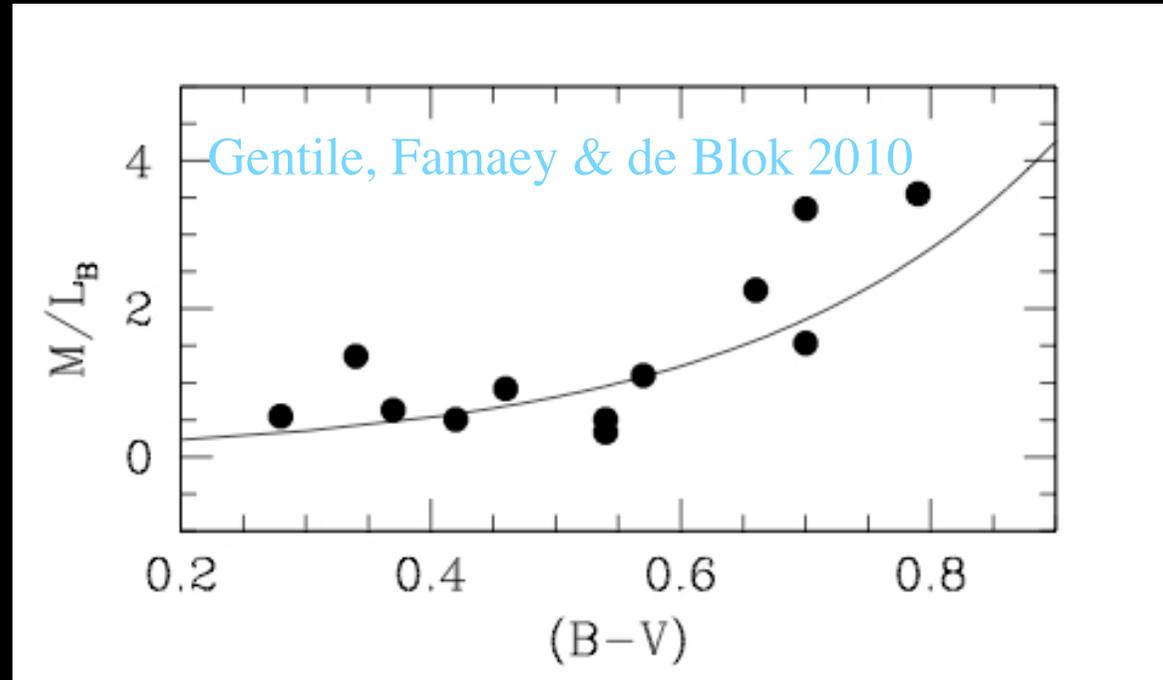
This formula fits >2000 galaxy rotation curves data points with *stellar M/L as the only free parameter* (*distance* can also slightly vary within observational errors)

$$\mu(x) = x / (1+x)$$



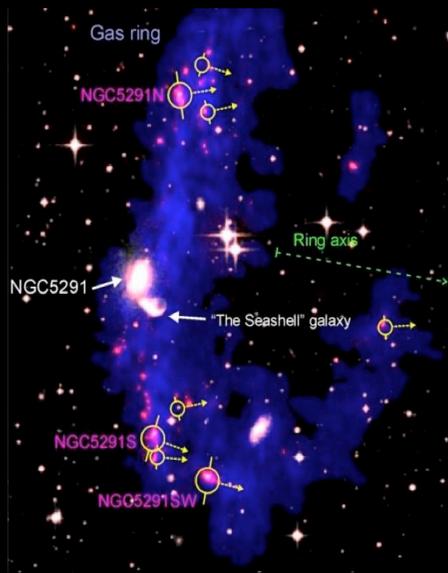
Might indicate a smaller value of $a_0 = 0.9 \times 10^{-8} \text{ cm/s}^2$

THINGS (Gentile, Famaey & de Blok 2010)



M/L follows predictions of population synthesis models

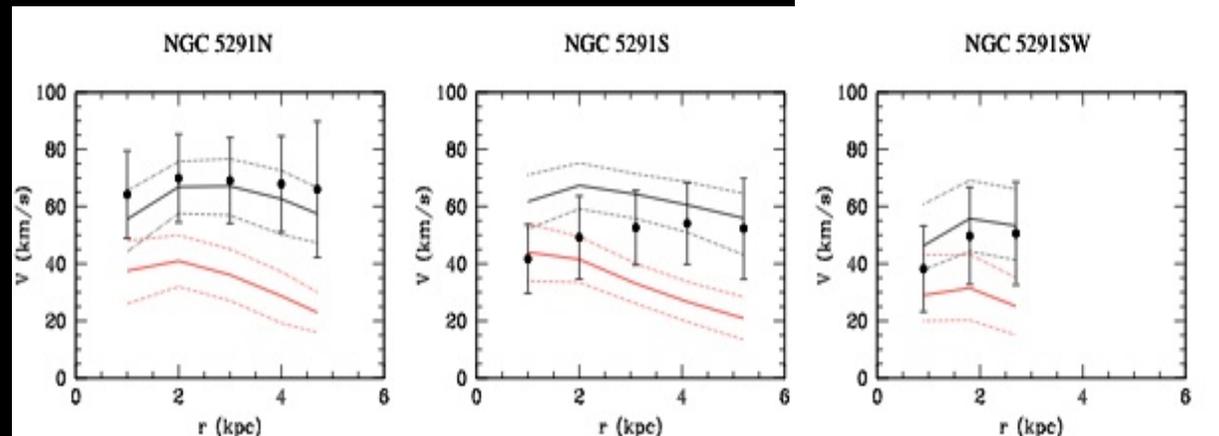
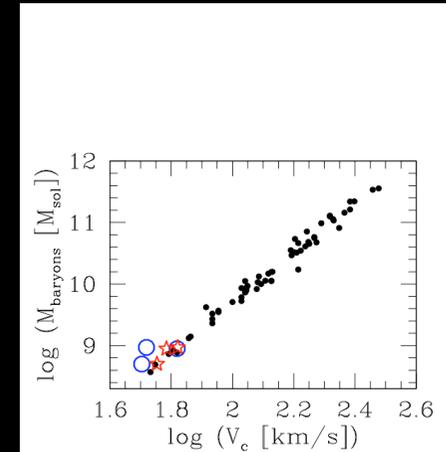
In BDM-dominated tidal dwarfs too



Tidal dwarf galaxies in
the NGC 5291 system

Bournaud et al. (2007)

Gentile, Famaey et al. (2007)



They fall on the BTF: why?

If Λ CDM is correct

- The success of the MOND formula in rotationally supported gas-dense galaxies is a consequence of the **feedback from the complex baryon physics** (stellar winds and supernova explosions)
- $\Lambda = (a_0/c)^2$ would just be a coincidence since DM and DE are completely unrelated sectors

\Rightarrow However, this all needs a true fine-tuning to reproduce this **history-independent** relation (both in CDM-dominated and BDM-dominated galaxies)

Formula cannot be exact

If we interpret it as *modified gravity*, for a point mass in the deep-MOND limit:

$$g = (GMa_0)^{1/2} / r$$

$$\Rightarrow F_m = m (GMa_0)^{1/2} / r$$

$$\text{BUT } F_M = M (Gma_0)^{1/2} / r$$

Corresponding modification of Newtonian gravity (MOND and QUMOND):

$$\nabla \cdot [\mu(|\nabla\Phi|/a_0) \nabla\Phi] = 4\pi G \rho_{\text{bar}}$$

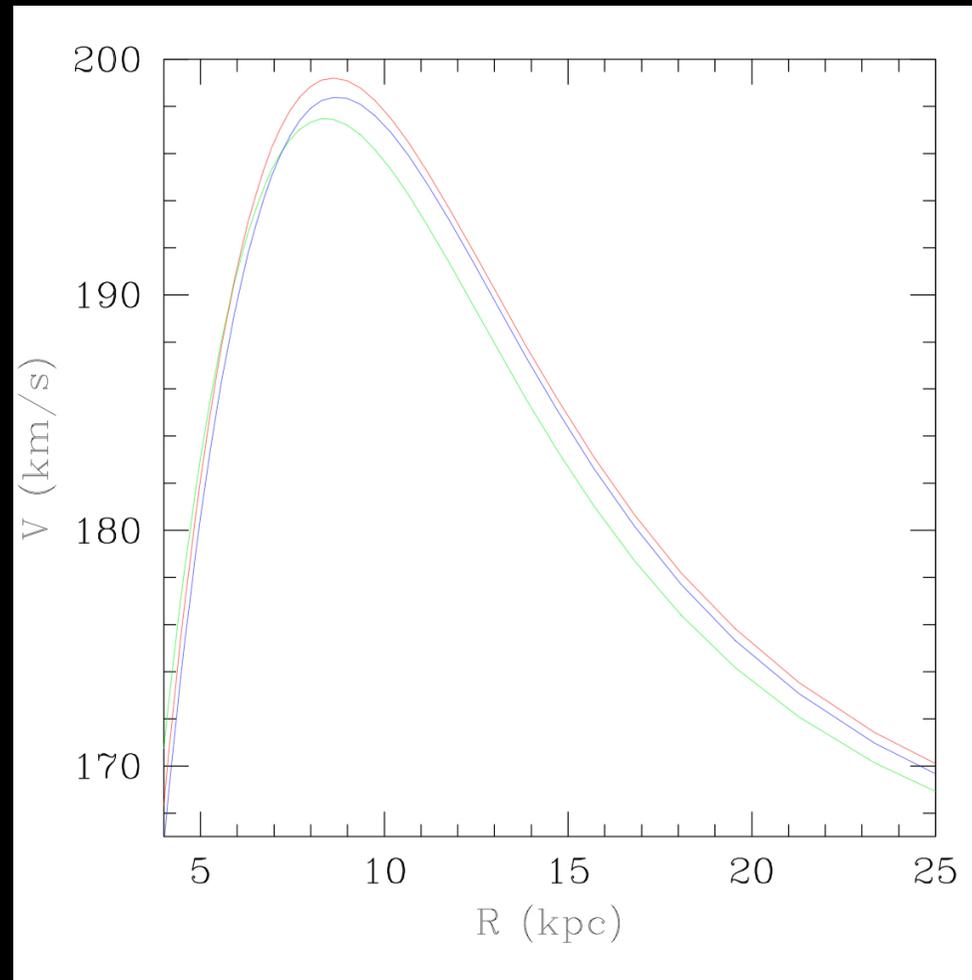
OR

$$\nabla^2 \Phi = \nabla \cdot [\nu(|\nabla\Phi_N|/a_0) \nabla\Phi_N]$$

Differing slightly
outside of spherical
symmetry

$$\Phi(r) \sim (GMa_0)^{1/2} \ln(r)$$

Comparing 3 formulations



Non-isolated systems

- In reality, *no* isolated systems: the external field in which an object is plunged influences the **internal** dynamics
- For instance, Milky Way in the slowly varying Great Attractor gravitational field ($0.01 a_0$)

- $\nabla \cdot [(\mathbf{g} + \mathbf{g}_e) \mu (|\mathbf{g} + \mathbf{g}_e| / a_0)] = \nabla \cdot (\mathbf{g}_n + \mathbf{g}_{ne})$

- In one dimension:

$$\mathbf{g}_n = \mathbf{g} \mu (|\mathbf{g} + \mathbf{g}_e| / a_0) + \mathbf{g}_e [\mu (|\mathbf{g} + \mathbf{g}_e| / a_0) - \mu (|\mathbf{g}_e| / a_0)]$$

When $|\mathbf{g}| \rightarrow 0$: $\mathbf{g}_n = \mathbf{g} \mu (|\mathbf{g}_e| / a_0)$, r^{-2} force, r^{-1} potential !

Escape speed

$$\frac{1}{2}v_{\text{esc}}^2(r) = \Phi(\infty) - \Phi(r)$$

Apply a $0.01a_0$ external field to the Milky Way, calculate the escape speed from the solar neighbourhood

-> $v_{\text{esc}} = 545 \text{ km/s}$ as observed !

Famaey, Bruneton & Zhao 2007

Wu et al. 2007

N-body simulations

- QUMOND simulations under development
- Classical MOND (Tiret & Combes) :
 - Bar appears more quickly than in CDM and are not slowed down by dynamical friction => fast bars
 - Long merging time-scale for interacting galaxies, but each oscillation could trigger a starburst
 - There must be *less* mergers at intermediate redshift than in CDM
 - It is possible to reproduce the Antennae, but with relatively fine-tuned initial conditions
 - It is possible to create warps in apparently isolated galaxies through the external field effect

Does MOND always work?

No: pressure-supported systems can be really problematic!

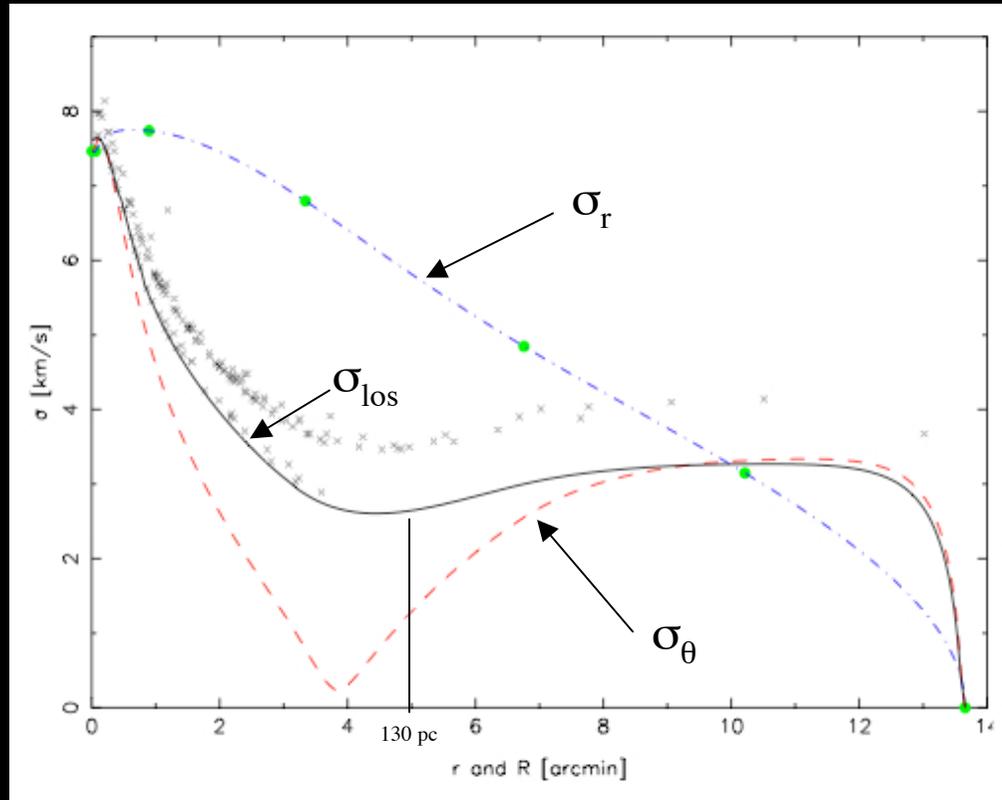
- **Galaxy clusters:** lensing and dynamics require additional dark matter (about as much as baryonic matter, a factor of 10 in the central parts)
- Velocity dispersion profiles and strong lensing of **elliptical galaxies:** generally **ok** in the field, but a few outliers inside groups and clusters
- Velocity dispersion profiles of **dwarf spheroidals:** generally **ok** but not (yet) for Sextans and Draco, and stability must be checked. The new ultra-faint dwarfs cannot be in equilibrium (**old TDGs?**)
- The total velocity dispersion in the **globular clusters Pal 14** and **Pal 4** (but not **Pal 3**) might be problematic for MOND (predicts 1 km/s instead of 0.5 km/s observed). But very few stars (**see Gentile et al. 2009**), orbit of the GC...?

Galaxy clusters

- **Ordinary neutrinos** of 2eV are **not enough** to explain the MOND discrepancy in X-ray groups
- Maybe a fermionic dark **HDM** particle? (**hot light sterile neutrinos with $m_\nu \sim 10\text{eV}$?**)
- **BUT note that** $\Omega_{\text{bvisible}} (=0.02) < \Omega_{\text{b}} (=0.04)$ at $z=0$
50% missing baryons \Rightarrow baryonic dark matter
How many baryons in WHIM?
- Bullet cluster \Rightarrow collisionless DM (Angus, Famaey, et al) \Rightarrow BDM in the form of e.g. **dense clumps of cold gas** (Pfenniger & Combes 1994), present only in galaxy clusters?
Microlensing? X-ray emission from cloud-cloud annihilation?
+ Why only in clusters and groups?

Globular clusters

NGC 2419 (Ibata et al.)



- Best MOND model 350 X less likely than best Newtonian model without dark matter

- BUT NGC 2419 part of the Virgo stream... on a very eccentric orbit => varying external field!

- **Are remote GCs generically on quite eccentric orbits? (cf Pal 14)**

A common solution for subgalactic and extragalactic scales?

- Modified MOND: **Bekenstein 2011** (see IAP talks 2011)
- Add a suitably chosen velocity scale s_0 such that s_0^2 can control the deepness of the potential and make a_0 effectively vary as :

$$a_{0\text{eff}} = a_0 \exp(-\Phi/s_0^2)$$

=> $a_{0\text{eff}}$ is **larger** in galaxy clusters (large $|\Phi|$) and **smaller** in globular clusters (small $|\Phi|$). Change in the zero-point is absorbed in definition of a_0

$$\nabla \cdot \left[\mu \left(\frac{|\nabla\Phi|}{a_{0\text{eff}}} \right) \nabla\Phi \right] + \frac{|\nabla\Phi|^2}{s_0^2} \mu \left(\frac{|\nabla\Phi|}{a_{0\text{eff}}} \right) - \frac{a_{0\text{eff}}^2}{s_0^2} F \left(\frac{|\nabla\Phi|^2}{a_{0\text{eff}}^2} \right) = 4\pi G\rho,$$

$$F = \int \mu \, dx$$

More predictions need cosmology, which needs a relativistic theory

- There now exists many different relativistic theories reproducing MOND
- The difficult thing is to have gravitational lensing and dynamics governed by the same potential
- In GR, the geodesic equation is:
$$d^2x^\mu / d\tau^2 = - \Gamma^\mu_{\alpha\beta} (dx^\alpha/d\tau) (dx^\beta/d\tau)$$
- reducing for timelike geodesics in weak-field to
$$d^2x^k / d\tau^2 = - \Gamma^k_{00} (dx^0/d\tau)^2 = - \Gamma^k_{00}$$

thus depending only on g_{00} (but **not** for null geodesics)

So $\Gamma^\mu_{\alpha\beta}$ plays the role of acceleration, but doesn't transform like a tensor => needs to introduce dark fields

Modifying GR

Einstein equations relate metric to stress-energy tensor just like Poisson equation relates potential to density. In **weak-field**:

$$g_{00} = - e^{2\Phi} = - (1 + 2\Phi)$$

$$g_{ij} = e^{2\Psi} \delta_{ij} = (1 + 2\Psi) \delta_{ij}$$

$$\Phi = -\Psi = \Phi_N \text{ in GR } (\Phi \Rightarrow \text{dynamics, } \Phi - \Psi \Rightarrow \text{lensing})$$

- **Idea:** replace GR with a theory reducing to the SAME metric but replacing Φ_N by Φ obeying MOND

- First try: add a scalar field and couple matter to

$$g'_{\alpha\beta} = e^{2\phi} g_{\alpha\beta}$$

with action of ϕ governed by a free function depending on $(\text{grad } \phi)^2$

\Rightarrow works for dynamics, but *doesn't work for lensing*

TeVes

Add a scalar field **and** a vector field and couple matter to:

$$g'_{\alpha\beta} = e^{-2\phi}(g_{\alpha\beta} + U_{\alpha}U_{\beta}) - e^{2\phi}U_{\alpha}U_{\beta}$$

with $g^{\alpha\beta} U_{\alpha}U_{\beta} = -1$, timelike in static situations,

and action $S = S_g + S_s + S_v + S_m$,

with action S_s of ϕ governed by a free function depending on $(\text{grad } \phi)^2$

$\Rightarrow \phi$ obeys a *B-M equation*, and plays the role of the dark matter potential (dynamics and lensing are governed by the **same** physical metric g')

In non-static situations, the spatial part of the vector field can play the role of DM (Ferreira & Starkman 2009 \Rightarrow DM in disguise), **can form structure but is most probably not enough to explain the missing mass in clusters nor the peaks of the CMB**

Other theories have slightly different static limits

Examples:

1) BIMOND (based on a function of a combination of $\Gamma^\mu_{\alpha\beta} - \Gamma'^\mu_{\alpha\beta}$ in the action):

$$\nabla^2 \Phi = \nabla \cdot [\nu (|\nabla\Phi_N| / a_0) \nabla\Phi_N]$$

2) Dipolar dark matter:

$$-\nabla \cdot [\mathbf{g} - 4\pi \mathbf{P}] = 4\pi G (\rho + \rho_{\text{dm}})$$

In any case...

- All these theories have difficulties in reproducing the CMB without dark matter (=> sterile neutrinos? Dipolar DM?)
- Independently from the theoretical framework, the MOND formula is an extremely efficient way of **predicting the gravitational field in rotationally supported galaxies**
- Any galaxy formation theory should be able to ultimately reproduce the MOND formula as a scaling relation for spirals (and TDGs)
- What makes it difficult is that it is history-independent!
- This is the goal of the complicated MOND theories that are developed and tested. Not getting rid of dark matter...